

STORAGE RING CROSS SECTION MEASUREMENTS FOR ELECTRON IMPACT SINGLE AND DOUBLE IONIZATION OF Fe¹³⁺ AND SINGLE IONIZATION OF Fe¹⁶⁺ AND Fe¹⁷⁺

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ABSTRACT

We report measurements of electron impact ionization for Fe¹³⁺, Fe¹⁶⁺, and Fe¹⁷⁺ over collision energies from below threshold to above 3000 eV. The ions were recirculated using an ion storage ring. Data were collected after a sufficiently long time that essentially all the ions had relaxed radiatively to their ground state. For single ionization of Fe¹³⁺, we find that previous single pass experiments are more than 40% larger than our results. Compared to our work, the theoretical cross section recommended by Arnaud & Raymond is more than 30% larger, while that of Dere is about 20% greater. Much of the discrepancy with Dere is due to the theory overestimating the contribution of excitation–autoionization via $n = 2$ excitations. Double ionization of Fe¹³⁺ is dominated by direct ionization of an inner shell electron accompanied by autoionization of a second electron. Our results for single ionization of Fe¹⁶⁺ and Fe¹⁷⁺ agree with theoretical calculations to within the experimental uncertainties.

Key words: atomic data – atomic processes – methods: laboratory

Online-only material: machine-readable tables

1. INTRODUCTION

Spectroscopic diagnostics of cosmic sources rely on accurate charge state distribution (CSD) calculations (Brickhouse 1996; Landi & Landini 1999). In stellar coronae, supernova remnants, galaxies, the intracluster medium of galaxy clusters, and other collisionally ionized plasmas, the balance between electron impact ionization (EII) and electron–ion recombination determines the CSD (Bryans et al. 2009). Thus, accurate EII data are needed to derive the CSD. For most objects, the temperature varies slowly enough that only electron impact single ionization (EISI) matters (Tendler et al. 1984), but when there is rapid heating, electron impact double ionization (EIDI) can be important (Müller 1986).

Most EII data come from theoretical calculations, as it is not possible to measure ionization for every ion. Experiments serve to benchmark these theoretical calculations. However, a major limitation of most existing EII measurements is that the ion beams used contained an unknown population of metastable ions. As the cross section for ionization from a metastable level generally differs from that for the ground level, the results of such experiments can be ambiguous and do not provide a clear test for theory.

Here we report EII measurements for Al-like Fe¹³⁺, Ne-like Fe¹⁶⁺, and F-like Fe¹⁷⁺. Of these three ion species, previous measurements exist only for Fe¹³⁺ (Gregory et al. 1987). However, those measurements were performed using a crossed-beams experiment and suffer from an unknown metastable fraction. We used an ion storage ring to avoid this problem. The ions were recirculated in the ring for several seconds before collecting data. This allowed any metastable levels to relax radiatively to the ground state. Our data thereby provide an unambiguous test for theoretical models.

The EISI cross section for Fe¹³⁺ was measured from about 300 eV to 3100 eV, where the direct ionization channels are

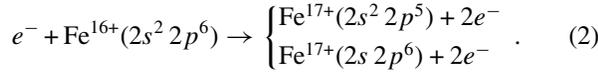
$$e^- + \text{Fe}^{13+}(2s^2 2p^6 3s^2 3p) \rightarrow \begin{cases} \text{Fe}^{14+}(2s^2 2p^6 3s^2) + 2e^- \\ \text{Fe}^{14+}(2s^2 2p^6 3s 3p) + 2e^- \\ \text{Fe}^{14+}(2s^2 2p^5 3s^2 3p) + 2e^- \\ \text{Fe}^{14+}(2s 2p^6 3s^2 3p) + 2e^- \end{cases} \quad (1)$$

The corresponding thresholds for 3p and 3s ionization are 392.2 eV and 421.2 eV, respectively (Kramida et al. 2012). Direct ionization of a 2p electron is possible above 1127.3 eV (Kramida et al. 2012) and ionization of the 2s is possible above about 1270 eV (Kaastra & Mewe 1993). However, following direct ionization of a principal quantum number $n = 2$ electron, the resulting excited state is expected to stabilize via autoionization with a probability of greater than 90%. This produces double ionization rather than single ionization (Kaastra & Mewe 1993). Ionization can also occur through excitation–autoionization (EA). For example, excitation of a 3s electron to a doubly excited level lying in the continuum is possible starting from the 392.16 eV ionization threshold. The system can then autoionize, resulting in EISI. EA via $n = 2$ to $n = 3$ excitations is predicted to occur starting at ~ 700 eV and via $2 \rightarrow 4 \sim 900$ eV (Pindzola et al. 1986; Arnaud & Raymond 1992; Dere 2007).

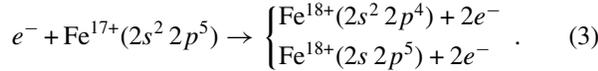
The EIDI cross section for Fe¹³⁺ forming Fe¹⁵⁺ was investigated from just below the threshold for direct double ionization at 848.4 eV up to 3100 eV. For highly charged ions the dominant double ionization process is expected to be single ionization of an electron in the $n = 2$ level forming a state that stabilizes by emission of a second electron (Müller & Frodl

1980). The threshold for this ionization–autoionization process is 1127.3 eV (Kramida et al. 2012).

The EISI cross sections for Fe^{16+} and Fe^{17+} were measured from 1100 eV to 3200 eV, where the direct ionization channels are



The ionization thresholds for these channels are 1262.7 eV for ionization of a $2p$ electron and 1394.7 eV for ionization of a $2s$ electron (Kramida et al. 2012). EA may occur through excitation of the $2s$ electron beginning from the ionization threshold of 1262.7 eV. Similarly, for Fe^{17+} the direct ionization channels are



These have ionization thresholds of 1357.8 for the $2p$ electron and 1472.2 for the $2s$ electron. Here again EA from excitation of a $2s$ electron is possible beginning at the ionization threshold of 1357.8 eV.

2. EXPERIMENTAL METHOD AND ANALYSIS

Cross section measurements were performed using the TSR heavy ion storage ring. This facility is located at the Max-Planck-Institut für Kernphysik in Heidelberg, Germany. The procedures used here are basically the same as those described by Linkemann et al. (1995) and Hahn et al. (2010, 2011a, 2011b, 2012a, 2012b). Below we outline the method and provide some details pertinent to the present measurements.

First a beam of iron ions was introduced into TSR. The ion beam energies were 141.0 MeV for Fe^{13+} , 176.9 MeV for Fe^{16+} , and 183.8 MeV for Fe^{17+} . In each case the isotope ^{56}Fe was used for the experiment. Two electron beams in the ring, dubbed the Cooler and the Target, were merged with the ions. Each electron beam is located in a different section of the ring. Initially the electron beams were used to cool the ion beam. That is, the energy of both electron beams was fixed to one where the electron velocity closely matched the average ion velocity, allowing elastic electron–ion collisions to reduce the energy spread of the ion beam (Poth 1990). This initial cooling period lasted three seconds.

During cooling, metastable states in the ion beam radiatively decayed. For Fe^{13+} there are two metastable levels with relatively long lifetimes. These are the $3s^2 3p^2 P_{3/2}$ level, whose decay to the ground state forms the well-known coronal green line (Edlén 1943; Esser et al. 1995), and the $3s 3p 3d^4 F_{9/2}$ level. The $3s^2 3p^2 P_{3/2}$ lifetime has been measured experimentally to be about 16.73 ms (Beiersdorfer et al. 2003; Brenner et al. 2007). The lifetime of the $3s 3p 3d^4 F_{9/2}$ level has been calculated theoretically to be 17.7 ms (Träbert et al. 1993). The similar lifetimes of these levels have been a source of systematic uncertainty in lifetime measurements at TSR (Träbert et al. 2002, 2009; Träbert 2010). For Fe^{16+} , the longest lived metastable levels are the $2s^2 2p^5 3s^3 P_{0,2}$, which have predicted lifetimes of 63.3 μs and 4.5 μs , respectively (Liang & Badnell 2010; Landi et al. 2012). The longest lived Fe^{17+} metastable level is the $2s^2 2p^5 2P_{1/2}$ level, which has a predicted lifetime of 51.5 μs (Del Zanna 2006; Landi et al. 2012). Since all these lifetimes are much shorter than the 3 s cooling time, the

metastable population for each ion is expected to be negligible during measurement.

After cooling, the Target was maintained at the cooling energy while the Cooler electron beam energy was varied so as to enable electron–ion collision studies at different energies. Ionized products from collisions in the Cooler section were diverted by a downstream dipole magnet onto a particle detector. The measurement energy was stepped through a range of energies. In between each measurement step, the ionization count rate was recorded for a fixed reference energy. This allowed us to assess the rate for background stripping off the residual gas. Ideally, this reference rate should be measured below the EII threshold, but at high measurement energies we were limited by the dynamic range of the electron beam power supply. Hence, for these energies the reference energy was set to a point where the ionization cross section was already measured in lower energy scans, allowing the background rate to be derived. The energy range for which the reference point was set below threshold was $E \leq 1050$ eV for Fe^{13+} EISI and $E \leq 1500$ eV for EIDI, $E \leq 1980$ eV for Fe^{16+} EISI, and $E \leq 1950$ eV for Fe^{17+} EISI.

The EII cross section σ_1 was obtained from the difference between the measured count rate and the background signal, normalized by the stored ion number and the electron density (Hahn et al. 2011a). The uncertainty introduced by the detector efficiency is about 3% (Rinn et al. 1982). Here and throughout all uncertainties are given at an estimated 1σ level. The electron density has an uncertainty of about 3% (Lestinsky et al. 2009). The stored ion number was derived from the ion current measured with a beam profile monitor (BPM; Hochadel et al. 1994). We calibrated the BPM several times during the measurement by comparing with the ion current measurement from a DC transformer (Unser 1981). The calibration was performed using currents of up to 21 μA (Hahn et al. 2011b). However, the DC transformer is not sensitive to the 0.5–5 μA currents present during measurement and could not be used directly for the analysis. We estimate that the uncertainty of the BPM contributes 15% to the experimental systematic uncertainty.

Energy-dependent pressure fluctuations change the background rate and can systematically distort the measured cross section. We corrected for these following Hahn et al. (2010). The magnitude of the correction was $(0.8 \pm 0.2)\%$ for Fe^{13+} EISI, $(4 \pm 2)\%$ for Fe^{16+} , and $(5 \pm 2)\%$ for Fe^{17+} . For the Fe^{13+} EIDI measurements, we did not find any systematic pressure fluctuations and so no correction was necessary. Because this correction could only be used when the reference point was below the threshold for ionization, there are other uncertainties on the cross sections in the higher energy ranges. The experimental uncertainties are given in Table 1.

3. RESULTS AND DISCUSSION

3.1. Single Ionization

3.1.1. Cross Sections

Figure 1 shows the EISI cross section for Fe^{13+} forming Fe^{14+} . These data are also available in the online edition of this journal as a table following the format of Table 2. In Figure 1 the filled circles show the measured cross section and the dotted curves illustrate the 1σ systematic uncertainty. The error bars on selected points represent the 1σ errors due to counting statistics. In some cases the error bars are smaller than the symbol size because the magnitude of the statistical uncertainty varies from

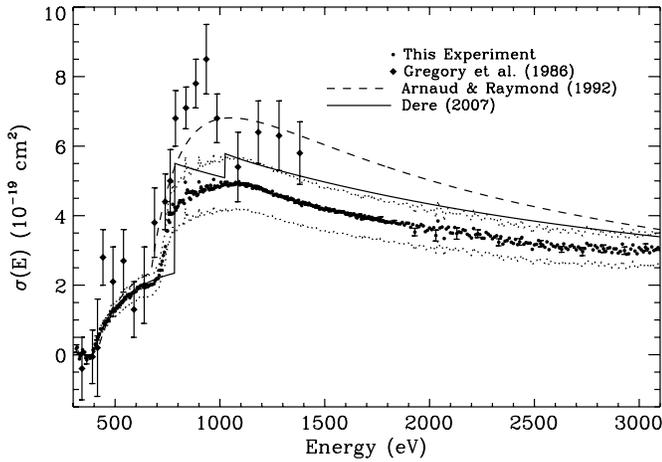


Figure 1. EISI cross section for Fe^{13+} forming Fe^{14+} (circles). The dotted curves illustrate the 1σ systematic uncertainties. Statistical uncertainties are indicated by the error bars on selected points, but in many cases they are smaller than the symbol size. The experimental results of Gregory et al. (1987) are shown by the diamonds. The dashed and solid curves show the theoretical cross sections given by Arnaud & Raymond (1992) and Dere (2007), respectively.

Table 1
Sources of Uncertainty

Source	Estimated 1σ Uncertainty			
	Fe^{13+}		Fe^{16+}	Fe^{17+}
	EISI	EIDI	EISI	EISI
Counting statistics	2%	3%	3%	6%
Detector efficiency	3%	3%	3%	3%
Ion current measurement	15%	15%	15%	15%
Electron density	3%	3%	3%	3%
Pressure fluctuations ^a	0.2% (0.8%)	...	2% (4%)	2% (5%)
Quadrature sum	16% (16%)	16%	16% (16%)	17% (17%)

Note. ^a The uncertainties in parentheses refer to the energy range where the reference point was above the ionization threshold. This is >1050 eV for Fe^{13+} EISI, >1500 eV for Fe^{13+} EIDI, >1980 eV for Fe^{16+} EISI, and >1950 eV for Fe^{17+} EISI. For Fe^{13+} EIDI no pressure fluctuations were observed.

Table 2
 Fe^{13+} Single Ionization Cross Section

E (eV)	σ_1 (cm^2)	Statistical Error (cm^2)
400	1.7490E-20	5.4474E-21
550	1.5980E-19	4.2844E-21
700	2.1955E-19	1.9516E-21
850	4.6733E-19	2.7333E-21
1000	4.8221E-19	2.4091E-21
1500	4.1659E-19	2.7346E-21
2000	3.6449E-19	1.6042E-20

Note. There is a systematic uncertainty of 16% in the cross section (see the text).

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about 1% to 4%, being smaller in places where more data were collected.

The diamonds in Figure 1 show crossed-beams EII measurements of Gregory et al. (1987). These results agree with our measurements from threshold to about 700 eV, but they

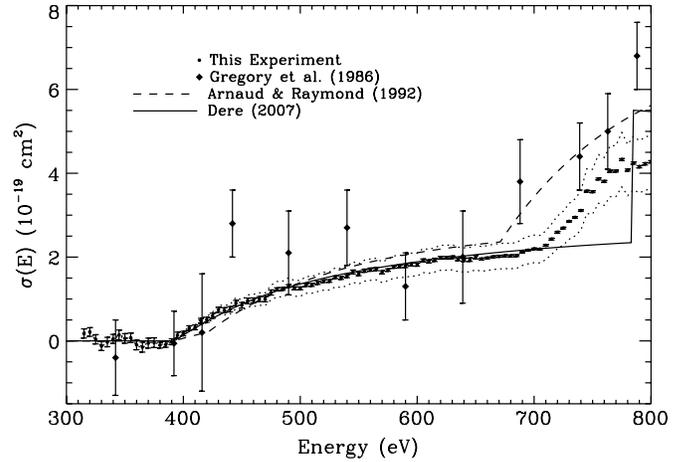


Figure 2. Portion of Figure 1 focusing on the threshold energy range.

are about 40% larger above 700 eV. This discrepancy is well outside the uncertainties of their and our measurements and is likely due to metastable ions in the crossed-beams experiment. Although Gregory et al. (1987) do not discuss the possible influence of metastables for this particular ion, the lifetimes of the Fe^{13+} metastable levels suggest that they could have been present. Their experiment used a 10 kV ion beam. Given that their device had a length scale of several meters, this implies that metastables with lifetimes $\gtrsim 10 \mu\text{s}$ could remain in the beam. As discussed above, the Fe^{13+} metastable levels $3s^2 3p^2 P_{3/2}$ and $3s 3p 3d^4 F_{9/2}$ have lifetimes of ≈ 16.7 ms and ≈ 17.7 ms, respectively, and could therefore both be present in the crossed beam experiment. The ionization threshold for the $3s 3p 3d^4 F_{9/2}$ level lies 81.86 eV below the ground-state ionization threshold (Kramida et al. 2012). However, the Gregory et al. (1987) cross section does not show any contribution below the ground-state threshold. This apparently low abundance of the $3s 3p 3d^4 F_{9/2}$ level may be due to its relatively high excitation energy and to not being strongly populated by cascades. It appears that it is primarily the $3s^2 3p^2 P_{3/2}$ level that is present in their beam, in addition to the ground level. This is also consistent with the observation that our results and those of Gregory et al. (1987) are in reasonable agreement at energies where direct ionization dominates, but disagree above the $2 \rightarrow 3$ EA threshold. The $3s^2 3p^2 P_{3/2}$ is part of the ground term and so it is expected to have a direct ionization cross section similar to that of the ground state, while the autoionization probabilities of levels above the $2 \rightarrow 3$ EA threshold could be different for the $^2 P_{3/2}$ state.

For comparison, Figure 1 illustrates the recommended cross section of Arnaud & Raymond (1992), which is based on the theoretical work of Younger (1983) and Pindzola et al. (1986). This cross section is up to 35% larger than our results. The reason for this is not clear. Also shown is the distorted wave results of Dere (2007), which agree with the measurement to within the level of the experimental uncertainties, though significant structural differences remain between the results of Dere (2007) and ours, as discussed below.

Excitation of a $3s$ electron to an autoionizing level is not included in these theoretical cross sections, but that appears not to be an issue here. Using the LANL Atomic Code (Magee et al. 1995), we have estimated that EA is possible for $3 \rightarrow n \gtrsim 10$ excitations and could contribute to the ionization cross section near threshold (Figure 2). Previous measurements for other $3s^2 3p^q$ ions have found a significant EA contribution via this

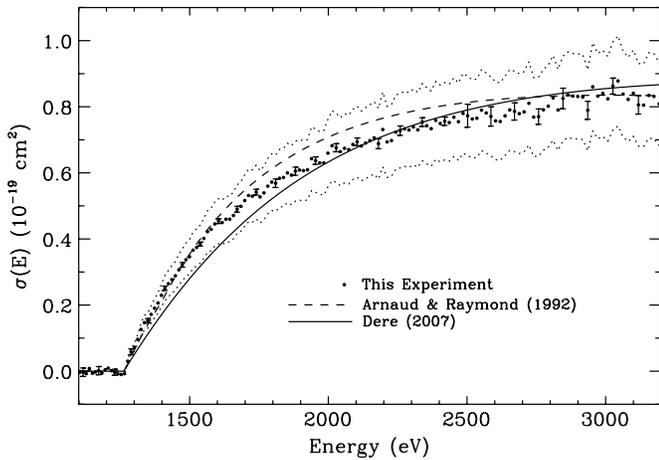


Figure 3. Same as Figure 1, but for Fe^{16+} forming Fe^{17+} .

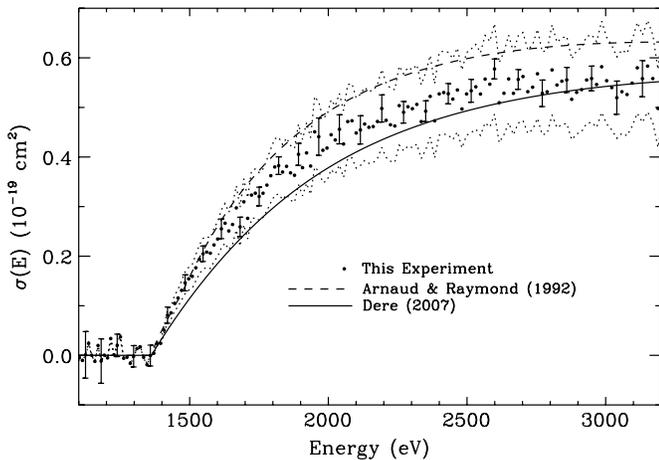


Figure 4. Same as Figure 3, but for Fe^{17+} forming Fe^{18+} .

channel ($q = 2$, Hahn et al. 2011b; $q = 3$, Hahn et al. 2011a; $q = 4$ and 5, Hahn et al. 2012a). This is supported by theoretical calculations ($q = 3$; Kwon & Savin 2012). However, here we find little discrepancy between theory and experiment near threshold, despite the omission of $3 \rightarrow n$ EA in the calculations. The n required for the excitation to autoionize is about the same for Fe^{13+} as for the ions of previous measurements, for example, for Fe^{11+} , $3 \rightarrow n > 8$ excitations were autoionizing. The reason this EA channel is relatively smaller here could be due to there being only one $3p$ electron. The effect could also be masked by the 16% systematic uncertainty.

At higher energies there are some discrepancies in the shape of the cross section. The Dere (2007) calculation overestimates the magnitude of $2 \rightarrow 4$ EA at about 1020 eV. This discrepancy has been found previously for similar ions (Hahn et al. 2011a, 2011b, 2012a). A possible explanation for the discrepancy is that the calculations underestimate the branching ratios for radiative stabilization. Another possibility is that the calculations underestimate the branching ratio for auto-double ionization (Kwon & Savin 2012), but this explanation is less likely since we do not observe any corresponding increase in the EIDI cross section at 1020 eV (see Section 3.2).

Figures 3 and 4 show the measured EISI cross sections for Fe^{16+} forming Fe^{17+} and Fe^{17+} forming Fe^{18+} , respectively. These data are available in the electronic edition of this journal as tables following the format of Tables 3 and 4. The figures

Table 3
 Fe^{16+} Single Ionization Cross Section

E (eV)	σ_1 (cm^2)	Statistical Error (cm^2)
1300.5	7.2389E-21	8.8629E-22
1499.3	3.4611E-20	6.0735E-22
1998.4	6.5945E-20	1.2343E-21
2502.2	7.7283E-20	3.4703E-21
3007.8	8.2407E-20	2.4911E-21

Note. There is a systematic uncertainty of 16% in the cross section (see the text).

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

Table 4
 Fe^{17+} Single Ionization Cross Section

E (eV)	σ_1 (cm^2)	Statistical Error (cm^2)
1407.7	5.0680E-21	1.6703E-21
1601.0	2.3464E-20	1.5598E-21
1995.0	4.1807E-20	3.7005E-21
2498.4	5.2723E-20	2.5672E-21
3003.6	5.5388E-20	2.5631E-21

Notes. There is a systematic uncertainty of 16% in the cross section (see the text).

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also illustrate the cross sections recommended by Arnaud & Raymond (1992), which is based on the calculations of Younger (1982), and the theoretical cross section of Dere (2007). For each ion there is generally good agreement with the measurement, to within the experimental uncertainties.

One discrepancy between experiment and theory is that the measured cross sections for both ions increase faster close to threshold than predicted by Dere (2007). This may be due to EA from excitation of a $2s$ electron to an energy above the threshold for ionization of the $2p$ electron. It should be noted, though, that the cross section of Arnaud & Raymond (1992) agrees very well with our results near threshold despite being based on calculations that included only direct ionization.

3.1.2. Rate Coefficients

Using the measured cross sections, we have derived EISI plasma ionization rate coefficients α_1 as a function of electron temperature T_e (cf. Hahn et al. 2011a). Figures 5, 6, and 7 show the results for Fe^{13+} , Fe^{16+} , and Fe^{17+} , respectively. In each case we compare these results to the rate coefficients of Arnaud & Raymond (1992) and Dere (2007). The vertical dotted lines in the figures indicate for ionization equilibrium the temperature ranges over which each ion is more than 1% abundant relative to the total Fe abundance as well as the temperature of peak abundance (Bryans et al. 2009).

Fe^{13+} is abundant from 1.3×10^6 K to 3.7×10^6 K with a peak abundance at 2.0×10^6 K. In this range the Fe^{13+} rate coefficients of Arnaud & Raymond (1992) differ from our experiment by up to 32%, while those of Dere (2007) agree with our results to within 10%. Fe^{16+} is abundant from 1.6×10^6 K to 1.3×10^7 K peaking at 4.1×10^6 K. Surprisingly, the older rate coefficients

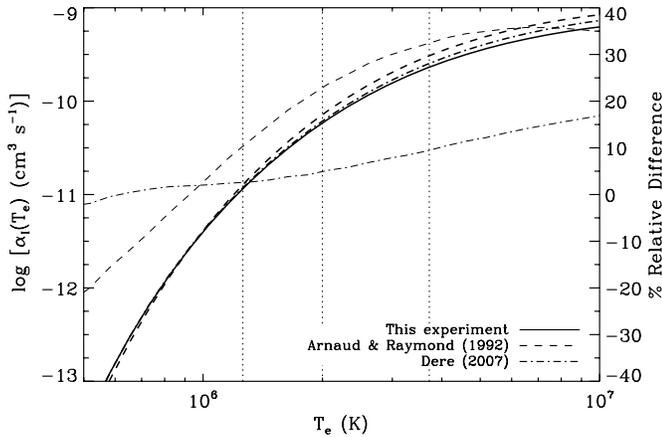


Figure 5. Thick lines show various plasma rate coefficients for Fe^{13+} forming Fe^{14+} . The solid curve indicates the experimental results, which can be read off the left axis. They are compared to the theoretical results of Arnaud & Raymond (1992, dashed curve) and Dere (2007, dash-dotted curve). The relative differences between these and the present results, (theory–experiment)/experiment, are shown by thin lines, with values read off the right axis. The dotted vertical lines denote the temperature range where Fe^{13+} is $>1\%$ abundant in collisional ionization equilibrium with the center line at the temperature of peak Fe^{13+} abundance (Bryans et al. 2009).

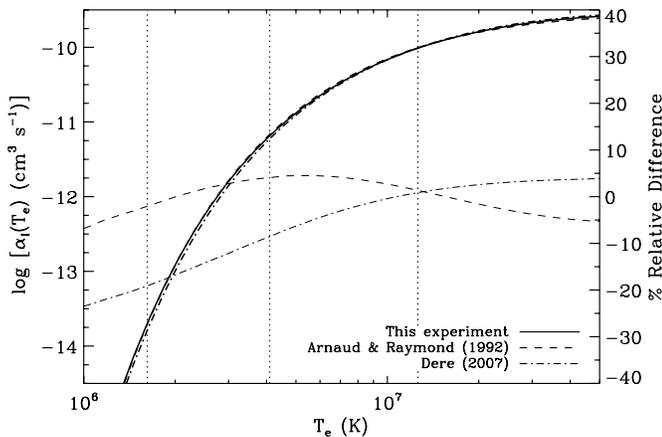


Figure 6. Same as Figure 5, but for Fe^{16+} forming Fe^{17+} .

of Arnaud & Raymond (1992) agree with our results to within 4% within this range. The more recent data of Dere (2007) differ by 19% at the low end of the temperature range due to the discrepancy near the ionization threshold. At the temperature of the peak abundance, they differ by 9% and at the high T_e end they agree with our measurement to within 1%. Finally, Fe^{17+} is abundant from 2.6×10^6 K to 1.5×10^7 K with peak abundance at 7.2×10^6 K. The rate coefficients of Arnaud & Raymond (1992) agree with our measurements in this range to within 13%. Those of Dere (2007) differ from the experimentally derived result by 17% at the low end of the T_e range (for the same reasons as for Fe^{16+}), by 9% at peak abundance, and by only 4% at the high T_e end.

The energy range covered by the experiment does not include excitation or ionization of a $1s$ electron. These channels are also neglected by Arnaud & Raymond (1992) and Dere (2007). Neglecting this may introduce a small error into the calculation of the ionization rate coefficient for Fe^{16+} and Fe^{17+} . The reason is that in deriving $\alpha_1(T_e)$, the integration over the cross section is performed up to $E_0 + 6k_B T_e$, where E_0 is the ionization threshold and k_B is the Boltzmann constant (Fogle et al. 2008). Thus, to

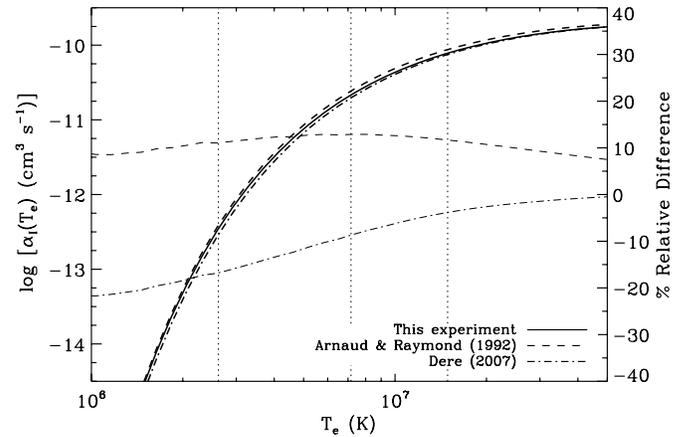


Figure 7. Same as Figure 5, but for Fe^{17+} forming Fe^{18+} .

calculate the rate coefficients over the full range where these ions are abundant, the integration should be performed up to 7984 eV for Fe^{16+} and up to 9113 eV for Fe^{17+} . These are greater than the measured energy range, and we have extrapolated the measurements to higher energies by scaling the Dere (2007) cross section to our measurements. This leaves an uncertainty as the integration limits exceed the $1s$ excitation and ionization thresholds. For Fe^{16+} the lowest EA channel is $1 \rightarrow 3$ EA, which opens at ≈ 7150 eV. For Fe^{17+} the lowest channel is $1 \rightarrow 2$ EA, which opens at ≈ 6450 eV (Hou et al. 2009; Kramida et al. 2012). The threshold for $1s$ ionization is 7714.7 eV for Fe^{16+} and 7823.2 eV for Fe^{17+} (Kramida et al. 2012).

In order to assess the possible error from neglecting direct ionization and EA, we have estimated the $1s$ ionization and excitation cross sections using the LANL Atomic Physics Code (Magee et al. 1995). These calculations show that the $1s$ direct ionization cross section is $\sim 0.5\%$ of the $n = 2$ direct ionization cross section for these ions in the relevant energy ranges. The maximum $1s$ EA cross section is the total excitation cross section. At the relevant energies, compared to the included EISI from $n = 2$, the Fe^{16+} $1 \rightarrow 3$ excitation cross section is $\sim 0.05\%$ and the Fe^{17+} $1 \rightarrow 2$ excitation cross section is about 0.5%. The contribution to single ionization from $n = 1$ excitation and ionization is actually smaller than implied by these cross sections. In the case of $1 \rightarrow n$ EA, the continuum state is estimated to radiatively stabilize 40% of the time and lead to EISI only for the other 60% (Kaastra & Mewe 1993). For $1s$ ionization, radiative relaxation of the intermediate state completes the EISI process 40%, but the system autoionizes the other 60%, leading to EIDI. Given the small cross sections for $1 \rightarrow n$ EA and $1s$ ionization, the branching ratios of the intermediate states, and the small fraction of the integrated energy range where they contribute at all, we expect that the omission of these processes has negligible effect on the calculated rate coefficients over the temperature ranges where Fe^{16+} and Fe^{17+} are abundant.

Table 5 presents coefficients for a polynomial fit to the scaled rate coefficient $\rho(x) = 10^{-6} \sum_i a_i x^i$, which can be used to reproduce the plasma rate coefficients. The rate coefficient $\alpha_1(T_e)$ is related to the scaled rate coefficient ρ by (Dere 2007)

$$\alpha_1(T_e) = t^{-1/2} E_0^{-3/2} E_1(1/t) \rho(x), \quad (4)$$

where $E_1(1/t)$ is the first exponential integral and $t = k_B T_e / E_0$ with E_0 the ionization threshold (392.2 eV for Fe^{13+} , 1262.7 for Fe^{16+} , and 1357.8 for Fe^{17+}). The scaled temperature x is

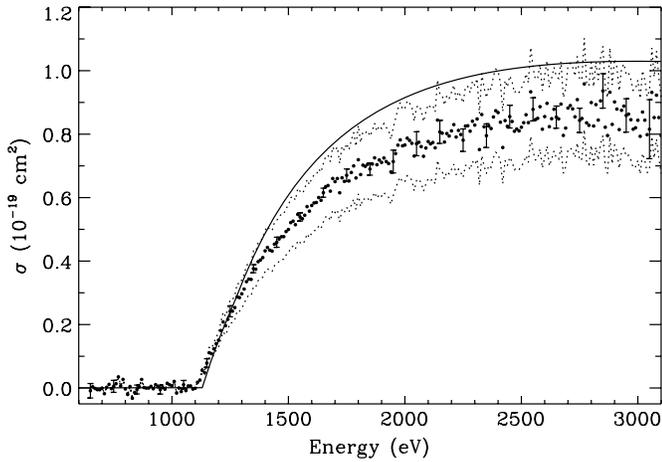


Figure 8. EIDI cross section for Fe^{13+} forming Fe^{15+} . The statistical uncertainty is indicated by error bars on selected points and the systematic uncertainties are illustrated by the dotted curves. The solid curve shows an estimate for the EIDI cross section due to direct ionization of an $n = 2$ electron forming a state that relaxes by emission of a second electron producing, in total, double ionization (see the text).

Table 5

Fifth-order Polynomial Fitting Parameters to Reproduce the Scaled Single Ionization Rate Coefficient $\rho = 10^{-6} \sum_{i=0}^5 a_i x^i \text{ cm}^3 \text{ s}^{-1} \text{ eV}^{3/2}$ (see Equations (4) and (6))

i	a_i		
	Fe^{13+}	Fe^{16+}	Fe^{17+}
0	9.59988	27.1411	17.4957
1	-53.2715	-24.4446	63.0820
2	401.820	43.0958	-485.718
3	-1018.29	74.8937	1517.84
4	1115.84	-252.597	-2127.02
5	-455.515	150.090	1094.75

given by

$$x = 1 - \frac{\ln 2}{\ln(t + 2)} \quad (5)$$

and by inverting T_e can be obtained from x :

$$T_e = \frac{E_0}{k_B} \left[\exp\left(\frac{\ln 2}{1-x}\right) - 2 \right]. \quad (6)$$

The experimental rate coefficients are reproduced to 1% accuracy or better for $T_e = 4 \times 10^5 - 1 \times 10^8$ K for Fe^{13+} , $T_e = 6 \times 10^5 - 1 \times 10^8$ K for Fe^{16+} , and $T_e = 1.5 \times 10^6 - 1 \times 10^8$ K for Fe^{17+} .

3.2. Double Ionization

Figure 8 shows the measured Fe^{13+} double ionization cross section. The dotted curves give the systematic uncertainty and the error bars illustrate the statistical uncertainty for select points. These data are available in the online version of this journal as a machine-readable version of Table 6. Although the threshold for double ionization is 848.4 eV, the cross section is consistent with zero until about 1100 eV when ionization of an $n = 2$ electron becomes possible. The solid line illustrates the expected double ionization cross section due to ionization of an $n = 2$ electron forming a state that relaxes through autoionization. This cross section was calculated using the LANL Atomic Physics Code (Magee et al. 1995) to determine

Table 6
 Fe^{13+} Double Ionization Cross Section

E (eV)	σ_1 (cm^2)	Statistical Error (cm^2)
900	5.5118E-24	1.3882E-21
1100	-7.6013E-23	1.2907E-21
1300	2.9696E-20	1.2073E-21
1500	4.9954E-20	1.2722E-21
2000	7.6579E-20	5.1443E-21
2500	8.2175E-20	3.5210E-21
3000	8.4374E-20	4.9785E-21

Note. There is a systematic uncertainty of 16% in the cross section (see the text).

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

the cross section for single ionization of an $n = 2$ electron and then scaling the result by the Auger yields of about 93% for $2p$ ionization and 95% for $2s$ ionization given by Kaastra & Mewe (1993). In the measured energy range, essentially all of the double ionization of Fe^{13+} is due to this process. This is consistent with results from other highly charged ions (Müller & Frodl 1980; Müller et al. 1985; Stenke et al. 1999; Hahn et al. 2011a, 2011b).

4. SUMMARY

We have measured cross sections for EISI from the ground states of Fe^{13+} , Fe^{16+} , and Fe^{17+} . For Fe^{13+} we find discrepancies of about 40% when compared to an earlier crossed-beams experiment. This is likely due to metastable ions in that work and their absence in ours. The theoretical cross section recommended by Arnaud & Raymond (1992) is more than 30% larger than our result. The recent calculation of Dere (2007) is also 20% larger than our result. The discrepancy with these theoretical calculations appears to be due to their treatment of EA. In particular, we do not observe the contribution from $2 \rightarrow 4$ EA predicted by Dere (2007). For Fe^{16+} and Fe^{17+} , our results generally agree with theory to within the experimental uncertainties. There is a small discrepancy in that the experimental cross section rises faster near threshold than predicted by Dere (2007). One possibility is that this is due to neglecting EA from the $n = 2$ excitations in the calculations. The measured EIDI cross section for Fe^{13+} is dominated by $n = 2$ ionization followed by autoionization of the excited state leading to a net double ionization. This result is consistent with measurements for other ions.

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